

Important Concepts in Electromagnetics

Coulomb's Law

Gauss's Law

Maxwell's equations

~~Static Boundary Conditions~~

Time Harmonic Maxwell's equations

Traveling waves

Helmholtz Equation

Time Harmonic Boundary Conditions

Transmission Lines

Gauss's Law + Divergence Theorem

$$\oint \vec{D} \cdot d\vec{S} = Q_{enc} \quad \text{Gauss's Law}$$

where $Q_{enc} = \int_V \rho_c dv$, or $\int_S \rho_c ds$, or $\int_V \rho_c dv$
Net electric flux $\Phi = Q$

Notes: $\oint \vec{D} \cdot d\vec{S} = |\vec{D}| \oint dS = Q_{enc} = \int_V \rho_c dv$

$$\Rightarrow |\vec{D}| = \dots \text{ done}$$

- this is how we solve for \vec{E} using
a Gaussian surface

$$\nabla \cdot \vec{D} = \rho_c \quad \text{pt form of Gauss}$$

Divergence Theorem:

$$\oint \vec{D} \cdot d\vec{S} = \int (\nabla \cdot \vec{D}) dv, \text{ from Gauss } \oint \vec{D} \cdot d\vec{S} = \int \rho_c dv = Q$$

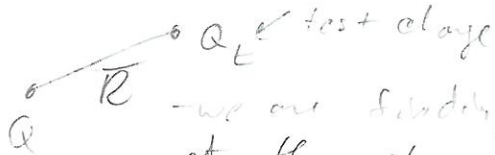
Coulomb's Law

- Coulomb's law is the force between 2 charges



$$F_1 = \frac{Q_1 Q_2}{4\pi \epsilon_0 |R_{21}|^2} \hat{r}_{21}$$

- Electric field from Coulomb's due to a pt. charge



- we are finding the E field from Q at the pt of the test charge

$$\vec{E} = \frac{\vec{F}_t}{Q_t} = \frac{1}{Q_t} \left[\frac{Q Q_t}{4\pi \epsilon_0 |R|^2} \hat{r}_1 \right]$$

$$\vec{E} = \frac{Q}{4\pi \epsilon_0 |R|^2} \hat{r}_1 \text{ due to a point charge}$$

- charge distributions

$$c = \frac{dQ}{dV} \left(\frac{C}{m^3} \right) \Rightarrow dQ = c dV \Rightarrow Q = \int c dV$$

$$d\vec{E} = \frac{dQ}{4\pi \epsilon_0 R^2} \hat{r} \Rightarrow \vec{E} = \int \frac{c dV}{4\pi \epsilon_0 R^2} \hat{r}$$

also works for line or surface

Displacement Current and Induced EMF

- from Ampere's law we know that

$$\Phi = \int \vec{B} \cdot d\vec{S} \quad \text{is the mag flux}$$

- due to displacement current ad induced EMF

$$V = -\frac{d\Phi}{dt}$$

- scalar potential

$$V = \int \vec{E} \cdot d\vec{e}$$

Laplace's Equation

- gives you voltage as a func of distance
- 1-D version

$$\frac{d^2V}{dy^2} = 0$$

- integrate both sides to get

$$V = Ay + B$$

- you can find

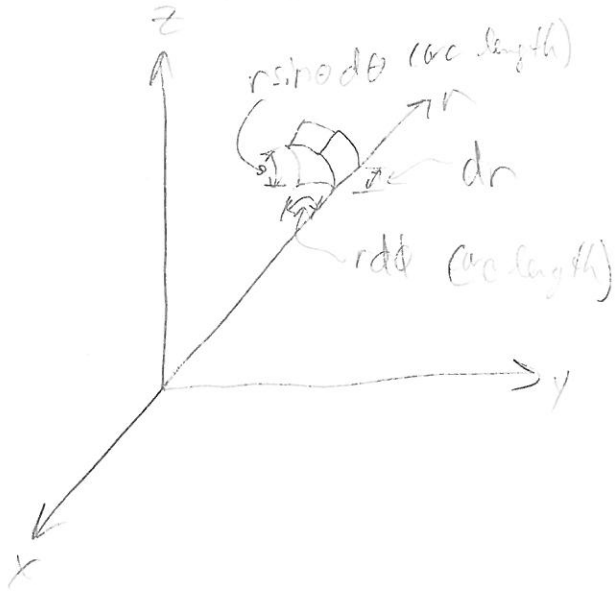
$$E = -\nabla V = -\nabla[Ay + B] \text{ to get solve for } A \text{ and } B$$

- or, if given voltage potentials at various locations, just subst. them into

$$V = Ay + B$$

A Note on Differential Components

- spherical coords



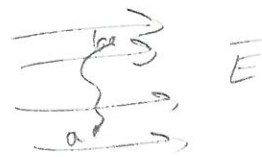
- you can make an $d\vec{S}_r$ comp's from this

$$\text{ie } d\vec{S}_r = \hat{r} r \sin \theta d\theta d\phi$$

- where \hat{r} is the unit normal vector to the surface

Electric Field Potential

$$V_{ab} = - \int_b^a \vec{E} \cdot d\vec{l} \quad (V)$$



Potential of a charge distribution:

$$V = \int_V \frac{e dv}{4\pi \epsilon_0 R}, \text{ because } dV = \frac{dQ}{4\pi \epsilon_0 R}$$

E-field and Potential

$$\vec{E} = -\nabla V$$

Potential due to a pt. charge

$$V = \frac{Q}{4\pi \epsilon_0 r}$$

Current Density

the continuity equation: $\nabla \cdot \vec{J} = -\frac{\partial \rho}{\partial t}$

current eq: $I = \int_S \vec{J} \cdot d\vec{S}$

Capacitance and Polarization

Polarization: a function of susceptibility χ_e

$$\bar{P} = \chi_e \epsilon_0 \bar{E}, \text{ where } \epsilon_r = 1 + \chi_e$$

$\Rightarrow \chi_e = \epsilon_r - 1$

$$\bar{D} = \epsilon_0 \epsilon_r \bar{E}$$

Capacitance: $C = \frac{Q}{V}$

Parallel platecap: $C = \frac{\epsilon_r \epsilon_0 A}{d}$

Static Boundary Conditions

Boundary conditions:

- 1) tangential component of \vec{E} is continuous across a dielectric layer

$$E_{t1} = E_{t2}, \quad \frac{D_{t1}}{\epsilon_{r1}} = \frac{D_{t2}}{\epsilon_{r2}}$$

- 2) normal component of \vec{D} is cont. across a dielectric layer

$$D_{n1} = D_{n2}, \quad \epsilon_{r1} E_{n1} = \epsilon_{r2} E_{n2}$$

- if a charge density, exists on boundary, then

$$D_{n1} - D_{n2} = -\rho_s, \quad \epsilon_{r1} E_{n1} - \epsilon_{r2} E_{n2} = -\frac{\rho_s}{\epsilon_0}$$

Ampere's Law and Magnetic Field

Biot-Savart's: $\vec{H} = \oint \frac{I d\vec{l} \times \vec{R}}{4\pi R^2}$

Ampere's Law: $I_{enc} = \oint \vec{H} \cdot d\vec{l}$

Relationship Between \vec{J} and \vec{H} : $\vec{J} = \nabla \times \vec{H}$

Magnetic Flux Density: $\vec{B} = \mu \vec{H}$

Magnetic Flux: $\Phi = \int_S \vec{B} \cdot d\vec{S}$

Magnetization: $\vec{B} = \mu_0 \vec{H} + \vec{M}$

$\vec{M} = \chi_m \mu_0 \vec{H}$

Vector Magnetic Potential: $\left\{ \begin{array}{l} \nabla \times \vec{A} = \vec{B} \\ \nabla \cdot \vec{A} = 0 \end{array} \right.$ (Must satisfy: (find \vec{B} first, then solve the diff eq for \vec{A}))

Current filament: $\vec{A} = \oint \frac{\mu I d\vec{l}}{4\pi R}$

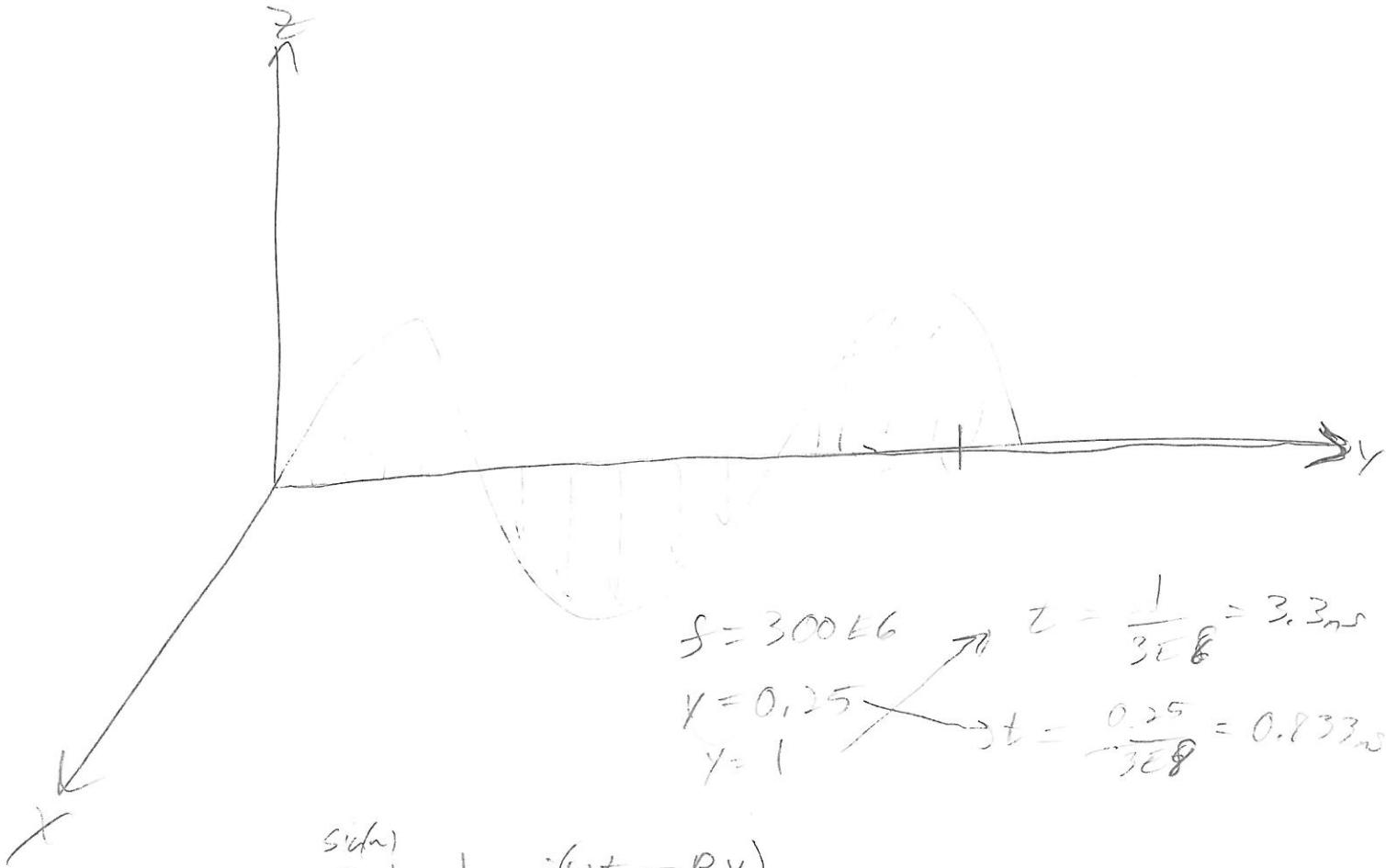
Sheet current: $\vec{A} = \int_S \frac{\mu \vec{K} d\vec{S}}{4\pi R}$

Volume current: $\vec{A} = \int_V \frac{\mu \vec{J} dV}{4\pi R}$

Stoke's Theorem: $\oint \vec{A} \cdot d\vec{l} = \int (\nabla \times \vec{A}) \cdot d\vec{S}$

- a way of finding the surface integral on a curl by using a single line integral

Electromagnetic waves



$$f = 300 \text{ EG}$$

$$v = 0,25$$

$$v = 1$$

$$z = \frac{1}{3 \text{ EG}} = 3,3 \text{ ns}$$

$$t = \frac{0,25}{3 \text{ EG}} = 0,833 \text{ ns}$$

$$\vec{E} = \hat{z} |E| e^{j(\omega t - \beta y)}$$

$$\cos(\omega t - \beta y) + j \sin(\omega t - \beta y)$$

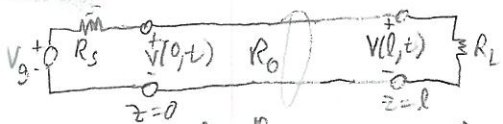
$$1,88 \text{ EG}$$

$$6,283$$

$$1 e^{j \left[\cancel{2\pi 300 \text{ EG}} (t) - j \frac{2\pi 300 \text{ EG}}{\cancel{3 \text{ EG}}} (y) \right]}$$

$$-13,1 \text{ E-3} + j 4,81$$

BOUNCE DIAGRAMS



$$\Gamma_L = \frac{R_L - R_0}{R_L + R_0} \quad \Gamma_S = \frac{R_S - R_0}{R_S + R_0}$$

$$T = \frac{L}{V} \text{ 1-way transit time}$$

$$0 < z < l \quad V(z,t) = \frac{R_0}{R_0 + R_S} V_g(t - z/v)$$

$$T < z < 2T \quad V(z,t) = \frac{R_0}{R_0 + R_S} \left[V_g(t - z/v) + \Gamma_L V_g(t - 2T + z/v) \right]$$

$$2T < z < 3T \quad V(z,t) = \frac{R_0}{R_0 + R_S} \left[V_g(t - z/v) + \Gamma_L V_g(t - 2T + z/v) + \Gamma_L \Gamma_S V_g(t - 4T + z/v) \right]$$

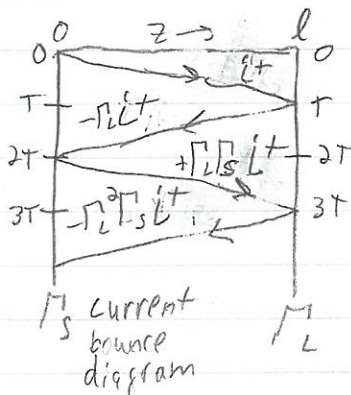
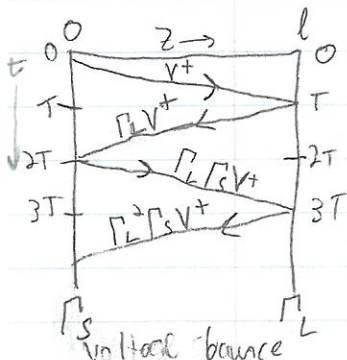
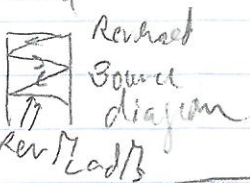
etc ...

for initially charged

T-lines, use

$$V^- = -V_0 R_0 / (R_L + R_0)$$

where



$$i^+ = \frac{V^+}{R_0}$$

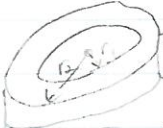
$$V^+ = \frac{V_0 R_0}{R_0 + R_S}$$

Inductance and Magnetic Circuits


Inductance: Flux linkage $\lambda = N\Phi$ where N = # turns
inductance $= L = \frac{\lambda}{I}$ where I is the coil current

Standard Conductor Configurations:


Toroid, square cross section:


$$L = \frac{\mu_0 N^2 a}{2\pi} \ln\left(\frac{r_2}{r_1}\right) \text{ (H)}$$

Toroid, general cross section S :


$$L \approx \frac{\mu_0 N^2 S}{2\pi} \text{ (H)} \quad (\text{assuming air flux density at average radius } r)$$


Parallel conductors of radius a :



radius a d l

$$\frac{L}{l} = \frac{\mu_0}{\pi} \cosh^{-1}\left(\frac{d}{2a}\right) \text{ (H/m)}, \quad \text{for } d \gg a, \quad \frac{L}{l} \approx \frac{\mu_0}{\pi} \ln \frac{d}{a}$$


cylindrical conductor parallel to ground plane:



radius a l

$$\frac{L}{l} = \frac{\mu_0}{2\pi} \cosh^{-1}\left(\frac{d}{2a}\right) \text{ (H/m)} \approx \frac{\mu_0}{2\pi} \ln\left(\frac{d}{a}\right) \text{ (H/m)}$$

long solenoid of small cross-sectional area S :


$$L = \frac{\mu_0 N^2 S}{l} \text{ (H)}$$

Time Harmonic Maxwell's Eqs

$$\nabla \times \vec{E} = -j\omega \vec{B} \quad \vec{B} = \mu \vec{H}$$

$$\nabla \times \vec{H} = j\omega \vec{D} + \vec{J}, \quad \vec{D} = \epsilon \vec{E} \text{ and } \text{with}$$

$$\vec{J} = \sigma \vec{E}$$

Maxwell's Equations

<u>Point Form:</u>	<u>Integral Form:</u>	<u>Name:</u>
$\nabla \times \vec{H} = \vec{J}_c + \frac{\partial \vec{D}}{\partial t}$	$\oint \vec{H} \cdot d\vec{l} = \int (\vec{J}_c + \frac{\partial \vec{D}}{\partial t}) \cdot d\vec{S}$	Ampere's Law
$\nabla \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$	$\oint \vec{E} \cdot d\vec{S} = \int (-\frac{\partial \vec{B}}{\partial t}) \cdot d\vec{S}$	Faraday's Law
$\nabla \cdot \vec{D} = \rho$	$\oint \vec{D} \cdot d\vec{S} = \int \rho dV$	Gauss's Law
$\nabla \cdot \vec{B} = 0$	$\oint \vec{B} \cdot d\vec{S} = 0$	Magnetic Monopoles do not exist